

# From invariant integrals to Haar weights

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# Introduction

For the [operator algebraic approach](#) to quantum groups, the existence of the Haar weights, the analogues of the Haar measures on a locally compact group, is crucial.

A first observation is that, up till now, there is [no general theory](#) of locally compact quantum groups where the existence of the Haar weights is a theorem, rather than part of the axioms.

On the other hand, [existence results](#) are available for [finite quantum groups](#), [compact quantum groups](#), [discrete quantum groups](#) and for [the dual](#) of a locally compact quantum group. Moreover Haar weights can be [constructed](#) following general ideas for most of the, also more complicated, [examples](#).

Therefore, having no general existence theorem is [not a big problem](#). Still, from a theoretical point of view, it would nice to have such a result.

In this talk, I will discuss [some aspects](#) of this. The results are available in the literature, but not so easy accessible, and often with more complicated arguments then necessary.

## Finite-dimensional Hopf algebras and duality

Let  $(A, \Delta)$  be a finite-dimensional Hopf algebra. Denote the dual by  $(B, \Delta)$  and use  $(a, b) \mapsto \langle a, b \rangle$  for the pairing. It is non-degenerate.

We consider  $A$  as a left  $A$ -module and as a left  $B$ -module:

$$(a, x) \mapsto ax \quad \text{and} \quad (b, x) \mapsto bx := \sum_{(x)} \langle x_{(2)}, b \rangle x_{(1)}$$

where  $a, x \in A$  and  $b \in B$ . Define  $V$  in  $B \otimes A$  by  $\langle V, a \otimes b \rangle = \langle a, b \rangle$ .

### Proposition

*For all  $x, x'$  we have  $V(x \otimes x') = \Delta(x)(1 \otimes x')$ .*

### Proof.

Let  $V = \sum_i b_i \otimes a_i$ . Then  $V(x \otimes x') = \sum_i b_i x \otimes a_i x'$  and so

$$V(x \otimes x') = \sum_{i, (x)} \langle x_{(2)}, b_i \rangle x_{(1)} \otimes a_i x' = \sum_{(x)} x_{(1)} \otimes x_{(2)} x'. \quad \square$$

We have  $\Delta(a) = V(a \otimes 1)V^{-1}$  and  $(S \otimes \iota)V = V^{-1}$ .

# Existence of a right integral

## Proposition

Let  $h : x \mapsto hx$  be a linear map from  $A$  to itself satisfying  $hb = S^2(b)h$ . Then  $\psi : a \mapsto \text{tr}(ah)$  is a right invariant functional.

## Proof.

Write  $V = \sum b_i \otimes a_i$ . Then for all  $a \in A$  we have

$$\begin{aligned}(\psi \otimes \iota)(V(a \otimes 1)V^{-1}) &= (\text{tr} \otimes \iota)(V(a \otimes 1)V^{-1}(h \otimes 1)) \\ &= \sum (\text{tr} \otimes \iota)((a \otimes 1)(1 \otimes a_i)V^{-1}(hb_i \otimes 1)) \\ &= \sum (\text{tr} \otimes \iota)((a \otimes 1)(1 \otimes a_i)V^{-1}(S^2(b_i)h \otimes 1))\end{aligned}$$

Now we have

$$(1 \otimes a_i)V^{-1}(S^2(b_i) \otimes 1) = (S \otimes \iota)(S(b_i) \otimes a_i)V.$$

We get 1 and this proves that  $(\psi \otimes \iota)\Delta(a) = \psi(a)1$ . □

We need enough such operators  $h$  to have that at least one such functional is non-zero.

### Proposition

For all linear maps  $y$  on  $A$  that commute with  $B$  we have that  $h := x \mapsto yS^{-2}(x)$  will satisfy  $hb = S^2(b)h$ .

### Proof.

$$\begin{aligned} S^{-2}(bx) &= \sum_{(x)} \langle x_{(2)}, b \rangle S^{-2}(x_{(1)}) \\ &= \sum_{(S^{-2}(x))} \langle S^2(S^{-2}(x))_{(2)}, b \rangle (S^{-2}x)_{(1)} \\ &= \sum_{(S^{-2}(x))} \langle (S^{-2}(x))_{(2)}, S^2(b) \rangle (S^{-2}x)_{(1)} \end{aligned}$$

and we see that

$$hbx = yS^{-2}(bx) = yS^2(b)S^{-2}(x) = S^2(b)yS^{-2}(x) = S^2(b)hx$$

□

One can now prove that there are elements  $y$  so that the corresponding  $\psi$  is non-zero. It is also true that a non-zero right invariant functional is automatically faithful.

# Uniqueness of integrals

## Proposition

Assume that  $\psi'$  is a right invariant functional and that  $\psi$  is a faithful right invariant functional. Then  $\psi'$  is a scalar multiple of  $\psi$ .

## Proof.

For all  $a, a' \in A$  we have

$$\Delta(a')(a \otimes 1) = \sum_{(a)} (\Delta(a'a_{(1)})(1 \otimes S(a_{(2)})).$$

If we apply  $\psi$  on the first factor, we obtain

$$(\psi \otimes \iota)(\Delta(a')(a \otimes 1)) = S((\psi \otimes \iota)((a' \otimes 1)\Delta(a)))$$

Now choose  $a$  so that  $\psi(\cdot a) = \varepsilon$ , the counit. Then we get for the left hand side  $a'$ .

Apply  $\psi'$  and use that  $\psi' \circ S$  is left invariant. Then we get  $\psi'(a') = \psi'(S(a))\psi(a')$ . □

## Existence for other cases

There are various other cases where existence theorems can be proven:

- for **compact** quantum groups,
- for **discrete** quantum groups,
- for the **dual quantum groups**.

The techniques vary and seem to be **far away** from the ones used to prove the existence of the Haar measure on a locally compact group.

As far as I know, there is **no general result** proving the existence of Haar weights on a general locally compact quantum group.

The proof that I gave for finite quantum groups might be a **basis** for a general existence result in the case where a multiplicative unitary (with the right properties) is already available. In fact, we have the **candidate** for the integral. The **missing point** is to show that it is non-trivial and densely defined.

The method has been used to find the Haar weights on the **examples** of locally compact quantum groups, with the multiplicative unitaries constructed by **Woronowicz**.



## Invariance of the Haar weights - Problem

There is a **major difficulty** proving invariance of the Haar weight on locally compact quantum groups.

In many cases, the invariant integral is first constructed on a **dense subalgebra**. Next it has to be **extended** to a weight on the operator algebra ( $C^*$ -algebra or von Neumann algebra). Finally one has to show that **invariance still holds** for this extension.

That this is not obvious is illustrated by the following result.

### Proposition

*Let  $\mathcal{H}$  be an infinite-dimensional Hilbert space and  $M$  the von Neumann algebra of all bounded linear operators on  $\mathcal{H}$ . Then there exists normal, semi-finite weights  $\varphi_1$  and  $\varphi_2$  on  $M$ , with orthogonal support, but equal on a weakly dense  $*$ -subalgebra.*

I will illustrate the solution to this problem for the **extension of the integral** on an algebraic quantum group to the generated locally compact quantum group. Observe that it is a problem that has to be solved in other situations as well.

# Algebraic quantum groups

We use the following terminology.

## Definition

A multiplier Hopf  $*$ -algebra with a positive right integral is called an **algebraic quantum group**.

Recall that a multiplier Hopf algebra is like a Hopf algebra but for algebras without identity. This means e.g. that the coproduct  $\Delta$  does not map into the algebraic tensor product but has values in the multiplier algebra  $M(A \otimes A)$ . It is assumed that elements of the form

$$\Delta(a)(1 \otimes a') \quad \text{and} \quad (1 \otimes a')\Delta(a)$$

belong to  $A \otimes A$ . Here we also assume that  $A$  is a  $*$ -algebra and that  $\Delta$  is a  $*$ -homomorphism.

A right integral is a linear functional  $\psi$  on  $A$  satisfying  $(\psi \otimes \iota)\Delta(a) = \psi(a)1$  in the multiplier algebra  $M(A)$ .

# Algebraic quantum groups - Duality

Denote by  $B$  the set of linear functionals on  $A$  of the form  $\psi(\cdot a)$  where  $a \in A$ .

## Theorem

*The space  $B$  is an algebraic quantum group where the product and the coproduct are adjoint to the coproduct and the product on  $A$ . The involution on  $B$  is defined by  $\langle a, b^* \rangle = \langle S(a)^*, b \rangle^-$ .*

We use  $(a, b) \mapsto \langle a, b \rangle$  on  $A \times B$ .

## Proposition

*There is an element  $V$  in the multiplier algebra  $M(B \otimes A)$  satisfying  $\langle V, a \otimes b \rangle = \langle a, b \rangle$  for all  $a, b$ .*

We need the extension of the pairing between  $B \otimes A$  and  $A \otimes B$  to  $M(B \otimes A) \times A \otimes B$ .

As in the finite-dimensional case,  $V$  is unitary and  $V^* = (S \otimes \iota)V$ .

# Algebraic quantum groups - Actions

Fix a positive right integral  $\psi$ . Denote by  $\mathcal{H}$  the Hilbert space completion of  $A$  with respect to the scalar product  $(x, x') \mapsto \psi(x'^* x)$ . We use  $x \mapsto \Lambda(x)$  for the canonical embedding of  $A$  in  $\mathcal{H}$ .

## Proposition

*There is a unitary operator on  $\mathcal{H} \otimes \mathcal{H}$ :*

$$\Lambda(x) \otimes \Lambda(x') \mapsto \sum_{(x)} \Lambda(x_{(1)}) \otimes \Lambda(x_{(2)}x').$$

One has the obvious left action of  $A$  on  $\Lambda(A)$  given by multiplication. We use  $a\Lambda(x) = \Lambda(ax)$ . We also have the left action of  $B$  on  $\Lambda(A)$  written as

$$b\Lambda(x) = \sum_{(x)} \langle x_{(2)}, b \rangle \Lambda x_{(1)}.$$

We combine this to an action of  $B \otimes A$  and extend it to the multiplier algebra  $M(B \otimes A)$ . Then we have that the unitary operator coincides with  $V$  on  $\Lambda(A) \otimes \Lambda(A)$ .

It is not entirely obvious, but as a consequence we get:

### Proposition

- i) *The algebra  $A$  acts with bounded operators on  $\mathcal{H}$ .*
- ii) *The algebra  $B$  acts with bounded operators on  $\mathcal{H}$ .*

Moreover we have  $*$ -representations and non-degenerate actions. In particular we have:

### Proposition

*For all  $b \in B$  we have  $\langle b\xi, \xi' \rangle = \langle \xi, b^*\xi' \rangle$ .*

### Proof.

Take  $x, x' \in A$  and  $b \in B$ . Then  $\langle b\Lambda(x), \Lambda(x') \rangle$  is equal to

$$\begin{aligned}\sum_{(x)} \langle x_{(2)}, b \rangle \psi(x'^* x_{(1)}) &= \sum_{(x')} \langle \mathcal{S}^{-1}(x'_{(2)}), b \rangle \psi(x'_{(1)} x) \\ &= \sum_{(x')} \langle x'_{(2)}, b^* \rangle^{-} \psi(x'_{(1)} x) \\ &= \langle \Lambda(x), b^* \Lambda(x') \rangle.\end{aligned}$$

□

# The associated locally compact quantum group

Define  $M$  as the von Neumann algebra generated by  $A$  on  $\mathcal{H}$  and a coproduct  $\Delta$  on  $M$  by  $\Delta(x) = V^*(x \otimes 1)V$ .

## Theorem

*The pair  $(M, \Delta)$  is a locally compact quantum group containing  $A$ . The right Haar weight on  $M$  is the unique extension of the right integral  $\psi$  on  $A$ .*

In fact, the right Haar weight is easy to define with the following notion from the theory of left Hilbert algebras:

## Definition

An element  $\eta \in \mathcal{H}$  is called *right bounded* if there is a bounded linear operator  $\pi_r(\eta)$  on  $\mathcal{H}$  satisfying  $\pi_r(\eta)\Lambda(a) = a\eta$  for all  $a \in A$ .

# The right Haar weight

## Definition

Define, for  $x \in M$  and  $x \geq 0$ ,

$$\bar{\psi}(x) = \sup\{\langle x\eta, \eta \rangle \mid \eta \text{ is right bounded and } \|\pi_r(\eta)\| \leq 1\}.$$

We will give some indications how it is shown that  $\bar{\psi}$  is a faithful normal semi-finite weight.

## Proposition

*The map  $\Lambda(a) \mapsto \Lambda(a^*)$  is pre-closed.*

We have

$$\langle \Lambda(a^*), \Lambda(a') \rangle = \psi(a'^* a^*) = \psi(a^* \sigma(a'^*)) = \langle \Lambda(\sigma(a'^*)), \Lambda(a) \rangle$$

This is used to prove (using left Hilbert algebra techniques) the following:

## Lemma

*The space of right bounded elements is dense.*

Again we need a property of the theory of left Hilbert algebras:

### Lemma

The set of right bounded elements  $\eta$  with  $\pi_r(\eta)^* = \pi_r(\eta')$  for some right bounded element  $\eta'$  gives a non-degenerate  $*$ -algebra of operators on  $\mathcal{H}$ .

### Proposition

If  $a \in A$  then  $\overline{\psi}(a^*a) = \psi(a^*a)$ .

### Proof.

For  $\eta$  right bounded we have

$$\langle a^*a\eta, \eta \rangle = \langle a\eta, a\eta \rangle = \langle \pi_r(\eta)\Lambda(a), \pi_r(\eta)\Lambda(a) \rangle = \langle \pi_r(\eta)^*\pi_r(\eta)\Lambda(a), \Lambda(a) \rangle.$$



One needs to prove additivity and then it follows that  $\overline{\psi}$  is a faithful normal semi-finite weight on  $M$  extending  $\psi$ .



## Right invariance of $\psi$

We now prove that this extension is still right invariant.

### Theorem

*The weight  $\bar{\psi}$  is right invariant.*

### Proof.

For all  $a \in A$  and  $\eta$  right bounded and any vector  $\xi$  we have

$$\Delta(a)(\eta \otimes \xi) = \sum_{(a)} (\pi_r(\eta) \otimes 1)(\Lambda(a_{(1)}) \otimes a_{(2)}\xi) = (\pi_r(\eta) \otimes 1)V(\Lambda(a) \otimes \xi).$$

By a result from the theory of left Hilbert algebras, this equation will also hold for  $x \in M$  satisfying  $\bar{\psi}(x^*x) < \infty$ . Then, when  $\bar{\psi}(x^*x) < \infty$ , we have

$$\bar{\psi}((\iota \otimes \langle \cdot, \xi, \xi \rangle)(\Delta(x^*x))) = \bar{\psi}(x^*x)\langle \xi, \xi \rangle$$

and  $\bar{\psi}$  is right invariant. □

# Conclusions

- We see that for the **passage** of algebraic quantum groups to locally compact quantum groups, some of the techniques used in the theory of **left Hilbert algebras** are needed, but not the full strength of this theory.
- On the other hand, for the development of the general theory of locally compact quantum group, you do need the **whole theory**.
- To understand locally compact quantum groups, a **good start** are the algebraic quantum groups. After all these contain already the **compact quantum groups** and their duals, the **discrete quantum groups**.
- We see that the proof of the existence of the Haar weights in the different situations uses **different techniques**. Understanding these different cases may help in the future to find a notion of locally compact quantum groups where the existence of the Haar weights is no longer part of the axioms, but a result, just like in the classical theory.

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